

Vapor–Liquid Equilibria, Density, and Speed of Sound for the System Poly(ethylene glycol) 400 + Methanol at Different Temperatures

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Solvent activity measurements have been performed on the system poly(ethylene glycol) 400 (PEG 400) + methanol at $T = (298.15 \text{ to } 328.15) \text{ K}$ using the improved isopiestic method. The obtained solvent activity data were correlated with the original and modified Flory–Huggins models. It was found that for the studied system at any mass fraction, as the temperature is increased the solvent activity and Flory parameter are both decreased, indicating an increase of interaction between polymer and methanol as the temperature is increased. The density and speed of sound data were also measured at $T = (288.15 \text{ to } 328.15) \text{ K}$ for the PEG 400 + methanol system. From the obtained density values, the excess molar volume data were calculated and fitted to the Redlich–Kister equation. The excess molar isentropic compression values calculated from the measured speed of sound and density data were also fitted to the Redlich–Kister equation.

Introduction

In recent years, numerous studies have been carried out on mixtures containing poly(ethylene glycols) (PEGs). PEG has numerous uses in biotechnology,^{1,2} in chemical partitioning,^{3–5} and in extractive crystallization of inorganic salts.⁶ An understanding of the thermodynamics of the polymer solutions is important in practical applications such as polymerizations, devolatilization, and the incorporation of plasticizers and other additives. Diffusion phenomenon in polymer melts and solutions are strongly affected by nonideal solution behavior, since the chemical potential rather than the concentration provides the driving force for diffusion. Proper design and engineering of many polymer processes depend greatly upon accurate modeling of thermodynamic parameters such as solvent activities.

The vapor–liquid data have been compiled by Wohlfarth⁷ for some polymer + solvent systems. Recently, in regard with PEG + alcohol solutions, a few vapor pressure data for the PEG + methanol and PEG + ethanol systems with polymer of molar mass 600 have been measured⁸ at 303.15 K by an apparatus based on the principle of electromicrobalance. In their vapor sorption measurements,⁸ however, they only considered the methanol and ethanol mass fractions in the limited ranges of 0.0143 to 0.3385 and 0.0298 to 0.3827, respectively. There are also vapor pressure or activity data for PEG + 2-propanol⁹ with different molar mass of PEG. However, as far as we know there is no vapor pressure or activity data for the PEG400 + methanol system at different temperatures in the literature. In this work, solvent activity data for PEG 400 + methanol were measured at $T = (298.15 \text{ to } 328.15) \text{ K}$ using the improved isopiestic method. The recent vapor pressure measurements on the solutions of LiBr in methanol¹⁰ provide us the osmotic coefficients of these solutions in wide ranges of concentrations and at different temperatures. These osmotic coefficient data permit us to use LiBr in methanol systems as an isopiestic standard; therefore, isopiestic measurement is now possible at higher concentrations and temperatures for solutions of other salts or polymer in methanol than was possible before.

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For PEG 400 + methanol, density and speed of sound data were also measured at $T = (288.15 \text{ to } 328.15) \text{ K}$, from which the values for excess molar volume and isentropic compression were calculated. For this system, there is no density or speed of sound data at different temperatures in the literature. The variation of activity, excess molar volume, and excess molar isentropic compression values with the polymer concentration and temperature provide us useful information in regard to polymer–solvent interactions and deviation of polymer solution from ideal solution.

Experimental Procedures

Materials. All the chemicals were obtained from Merck. Lithium bromide with purity of minimum mass fraction 0.995 was dried in an electrical oven at about 110 °C for 24 h prior to use. Methanol with purity of minimum mass fraction 0.995 was dehydrated according to Vogel.¹¹ Double-distilled, deionized water was used. For molar mass of PEG, the value 400 $\text{g}\cdot\text{mol}^{-1}$ was used.

Apparatus and Procedures. The isopiestic apparatus employed is essentially similar to the one used previously.¹² Recently this technique has been used for the measurement of activity of the methanol in alcohol + poly(vinyl pyrrolidone),¹³ alcohol + poly(propylene glycol),¹⁴ and 2-propanol in 2-propanol + poly(ethylene glycol)⁹ systems with different molar masses of the polymer. This apparatus consisted of a five-leg manifold attached to round-bottom flasks. Two flasks contained the standard LiBr solutions, two flasks contained PEG solutions, and the central flask was used as a methanol reservoir. The apparatus was held in a constant-temperature bath at least 120 h for equilibration at $T = (298.15 \text{ to } 328.18) \text{ K}$. The temperature was controlled to within $\pm 0.005 \text{ }^\circ\text{C}$ by a Heto temperature controller (Heto therm PF, Heto Lab Equipment, Denmark). After equilibrium had been reached, the manifold assembly was removed from the bath, and each flask was weighed with an analytical balance (Shimatzu, 321-34553, Shimatzu Co., Japan) with an uncertainty of $\pm 1 \times 10^{-7} \text{ kg}$. It was assumed that the equilibrium condition was reached when the differences between the mass fractions of each duplicate were less than 1 %. In all

Table 1. Experimental Isopiestic Mass Fraction w , Osmotic Coefficients Φ , and Methanol Activity a_1 for Methanol (1) + PEG400 (2) at Different Temperatures

w_{LiBr}	w_2	Φ_{LiBr}	a_1	p/kPa	w_{LiBr}	w_2	Φ_{LiBr}	a_1	p/kPa	w_{LiBr}	w_2	Φ_{LiBr}	a_1	p/kPa
$T = 298.15 \text{ K}$														
0.0148	0.1076	0.909	0.9900	16.784	0.0963	0.4389	1.017	0.9232	15.637	0.2168	0.7307	1.665	0.7116	12.019
0.0190	0.1240	0.905	0.9871	16.734	0.1159	0.4993	1.087	0.9001	15.241	0.2395	0.7732	1.844	0.6514	10.993
0.0311	0.1547	0.900	0.9789	16.593	0.1260	0.5276	1.128	0.8869	15.015	0.2569	0.8033	1.995	0.6012	10.139
0.0411	0.2298	0.904	0.9718	16.472	0.1786	0.6538	1.403	0.7985	13.502	0.2875	0.8570	2.294	0.5050	8.506
0.0580	0.3028	0.924	0.9588	16.248	0.1836	0.6688	1.435	0.7881	13.325	0.2966	0.8643	2.391	0.4752	8.001
0.0777	0.3770	0.964	0.9418	15.957	0.1936	0.6867	1.499	0.7669	12.962					
$T = 308.15 \text{ K}$														
0.0146	0.0846	0.892	0.9903	27.684	0.1131	0.4867	1.075	0.9038	25.228	0.2167	0.7226	1.646	0.7146	19.881
0.0307	0.1664	0.876	0.9798	27.386	0.1335	0.5402	1.162	0.8763	24.449	0.2439	0.7788	1.854	0.6433	17.876
0.0431	0.2354	0.883	0.9711	27.138	0.1477	0.5830	1.230	0.8545	23.831	0.2643	0.8102	2.029	0.5839	16.208
0.0587	0.3008	0.908	0.9591	26.797	0.1638	0.6173	1.314	0.8271	23.056	0.2898	0.8515	2.273	0.5045	13.985
0.0772	0.3646	0.953	0.9429	26.337	0.1702	0.6353	1.349	0.8154	22.725	0.3218	0.8926	2.626	0.3987	11.032
0.0976	0.4381	1.017	0.9221	25.747	0.1886	0.6773	1.459	0.7785	21.683	0.3492	0.9214	2.976	0.3078	8.504
$T = 318.15 \text{ K}$														
0.0370	0.1948	0.918	0.9743	43.380	0.1348	0.5358	1.156	0.8756	38.901	0.2654	0.8039	2.013	0.5848	25.817
0.0462	0.2430	0.920	0.9676	43.076	0.1488	0.5739	1.220	0.8543	37.937	0.2893	0.8392	2.234	0.5111	22.528
0.0616	0.3137	0.936	0.9557	42.535	0.1653	0.6098	1.305	0.8264	36.676	0.3154	0.8762	2.502	0.4273	18.800
0.0770	0.3682	0.964	0.9424	41.931	0.2042	0.6934	1.540	0.7471	33.099	0.3254	0.8876	2.614	0.3944	17.340
0.1054	0.4623	1.042	0.9134	40.614	0.2105	0.7094	1.583	0.7323	32.433					
0.1172	0.4888	1.084	0.8992	39.970	0.2517	0.7793	1.895	0.6248	27.607					
$T = 328.15 \text{ K}$														
0.0198	0.1117	0.714	0.9848	22.686	0.0664	0.3313	0.806	0.9409	21.659	0.1607	0.6431	1.314	0.7658	17.577
0.0285	0.1688	0.724	0.9777	22.520	0.0983	0.4436	0.920	0.8991	20.682	0.1675	0.6618	1.375	0.7455	17.105
0.0304	0.1776	0.727	0.9761	22.482	0.1157	0.5042	1.005	0.8698	19.998	0.1825	0.7086	1.522	0.6973	15.987
0.0557	0.2859	0.777	0.9525	21.930	0.1446	0.5975	1.187	0.8082	18.563					

cases, averages of the duplicate are reported as the total isopiestic mass fraction. The uncertainty in the measurement of solvent activity was estimated to be ± 0.0002 .

In this study, the ultrasonic velocity and density of PEG 400 + methanol mixtures were also measured at different temperatures with a digital vibrating-tube analyzer (Anton Paar DSA 5000, Austria) with a proportional temperature controller that kept the samples at working temperature with an uncertainty of 0.001 K. Recently this apparatus has been used for measurement of density and sound velocity of aqueous solutions of poly(propylene glycol).¹⁵ The apparatus was calibrated at each temperature with distilled water and dry air. The apparatus was also tested with the density of a known molality of aqueous NaCl using the data of Pitzer et al.¹⁶ The uncertainty of the instrument is $\pm 0.003 \text{ kg}\cdot\text{m}^{-3}$ for density and $0.1 \text{ m}\cdot\text{s}^{-1}$ for ultrasonic velocity.

Results and Discussion

Experimental Results. At isopiestic equilibrium, the activity of solvent in the reference and PEG solutions must be the same. Therefore, the isopiestic equilibrium mass fractions with reference standard solutions of LiBr in methanol as reported in Table 1, enabled the calculation of the solvent activity, a_1 , in the solutions of methanol (1) + PEG (2) from that of reference solutions using the relation

$$\ln a_1 = -\nu m_{\text{LiBr}} \phi_{\text{LiBr}} M_s, \quad m_{\text{LiBr}} = \frac{w_{\text{LiBr}}}{M_{\text{LiBr}}(1 - w_{\text{LiBr}})} \quad (1a,b)$$

where ν is the sum of stoichiometric numbers of anion and the cation in the reference solutions; m_{LiBr} and w_{LiBr} are respectively the substance concentration and mass fraction of LiBr, which is in isopiestic equilibrium with the polymer solutions; M_s and M_{LiBr} are respectively the molar masses of the solvent and LiBr; and ϕ_{LiBr} is the osmotic coefficient of the isopiestic reference standard, calculated at m_{LiBr} . The necessary ϕ_{LiBr} values at any

m_{LiBr} were obtained from the fitted extended Pitzer ion interaction model of Archer equation:

$$\phi_{\text{LiBr}} = \left[1 - A_\phi \left(\frac{\sqrt{m_{\text{LiBr}}}}{1 + b\sqrt{m_{\text{LiBr}}}} \right) \right] + m_{\text{LiBr}}(\beta^{(0)} + \beta^{(1)} \exp[-\alpha_1 \sqrt{m_{\text{LiBr}}}] + \beta^{(2)} \exp[-\alpha_2 \sqrt{m_{\text{LiBr}}}] + (m_{\text{LiBr}})^2 (C^{(0)} + C^{(1)} \exp[-\alpha_3 \sqrt{I}]) \quad (2)$$

The parameters of eq 2 for LiBr in methanol solutions in the concentration range (0.0411 to 6.8675) $\text{mol}\cdot\text{kg}^{-1}$ and temperature range of $T = (298.15 \text{ to } 333.15) \text{ K}$ have been reported by Nasirzadeh et al.¹⁰ It was also shown that,¹⁰ using $\alpha_1 = 2$, $\alpha_2 = 7$, $\alpha_3 = 1$, and $b = 3.2$, the osmotic coefficients ϕ_{LiBr} are reproducible with standard deviation of about 0.01 in the above concentration and temperature ranges. The obtained methanol activity data for the system methanol (1) + PEG400 (2) are given in Table 1. To see the effect of temperature on the methanol activity, the a_1 values were plotted versus polymer mass fractions at (298.15 and 328.15) K in Figure 1. As can be seen from Figure 1, for a given polymer mass fraction an increase in temperature causes a slight decrease in its activity value.

From the calculated solvent activity data, the vapor pressure of methanol solutions, p , were determined with the help of the following relation:

$$\ln(a_1) = \ln\left(\frac{p}{p^*}\right) + \frac{(B - V_s^*)(p - p^*)}{RT} \quad (3)$$

where B , V_s^* , and p^* are the second virial coefficient, molar volume, and vapor pressure of pure methanol, respectively. The values of the physical properties for the methanol, which are taken from Nasirzadeh et al.,¹⁰ are summarized in Table 2. The experimental vapor pressure data are also given in Table 1.

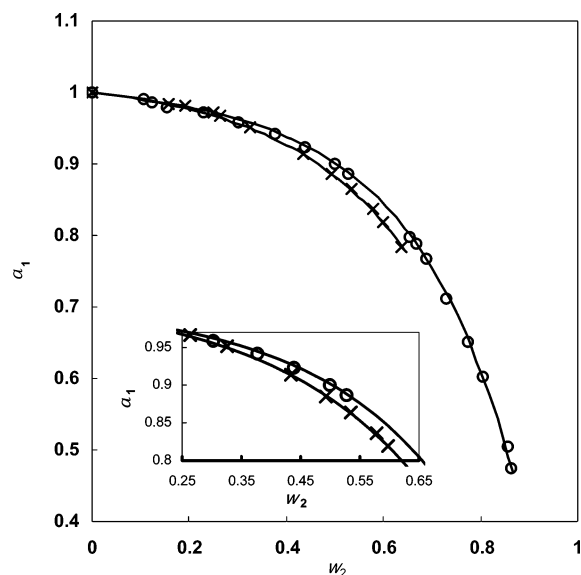


Figure 1. Experimental and calculated activity of methanol, a_1 , plotted against mass fraction of polymer, w_2 , for methanol (1) + PEG400 (2) system at two temperatures: \circ , 298.15 K; \times , 328.15 K; $-$, modified Flory–Huggins model (eq 7).

Table 2. Physical Properties of Methanol^a

T K	$10^5 V_s^*$ $\text{m}^3 \cdot \text{mol}^{-1}$	$10^3 B$ $\text{m}^3 \cdot \text{mol}^{-1}$	ρ $\text{kg} \cdot \text{m}^{-3}$	p^* kPa
298.15	4.075	-1.9229	786.373	16.956
308.15	4.124	-1.5317	776.894	27.960
318.15	4.176	-1.2410	767.303	44.550
328.15	4.230	-1.0222	757.573	68.791

^a Taken from ref 10.

Experimental data on density ρ and ultrasonic velocity u of various methanol + PEG solutions determined at $T = (288.15$ to $328.15)$ K are given in Table 3.

The excess molar volumes, V_m^{ex} , and the isentropic compression, $K_{s,m}^{\text{ex}} = -(\partial V_m^{\text{ex}}/\partial P)_s$ were determined by the following expressions:

$$V_m^{\text{ex}} = \sum_{i=1}^2 x_i M_i \left(\frac{1}{\rho} - \frac{1}{\rho_i} \right) \quad (4)$$

$$K_{s,m}^{\text{ex}} = \sum_{i=1}^2 x_i M_i \left[\frac{1}{(\rho u)^2} - \frac{1}{(\rho_i u_i)^2} \right] \quad (5)$$

where x is the mole fraction; M is the molar mass; and subscripts 1 and 2 stand for methanol and polymer, respectively. The values obtained for V_m^{ex} and $K_{s,m}^{\text{ex}}$ at different concentrations and temperatures are also reported in Table 3. A plot of V_m^{ex} and $K_{s,m}^{\text{ex}}$ values versus polymer mole fraction are shown respectively in Figures 2 and 3. These figures show that at different concentrations and working temperatures both the V_m^{ex} and $K_{s,m}^{\text{ex}}$ are negative and become more negative when temperature increases. The negative V_m^{ex} value can be explained in terms of positive contributions due to breaking of like interactions of the pure liquids and negative contributions due to the formation of unlike (polymer–solvent) interactions and also accommodation of small methanol molecules in the voids provided by the PEG (packing effect). The interaction between the hydrogen

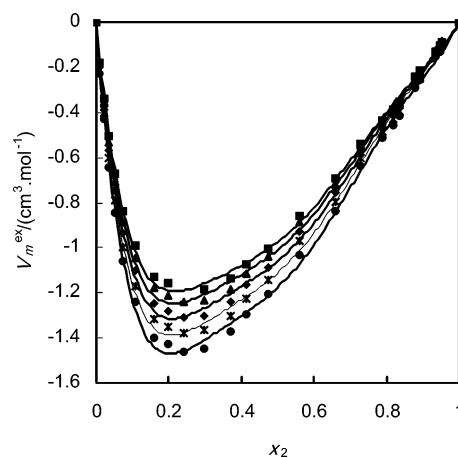


Figure 2. Plot of excess molar volume $V_m^{\text{ex}} (\text{cm}^3 \cdot \text{mol}^{-1})$ for methanol (1) + PEG400 (2) system against mole fraction of polymer x_2 at different temperatures: \blacksquare , 288.15 K; \blacktriangle , 298.15 K; \blacklozenge , 308.15 K; \times , 318.15 K; \bullet , 328.15 K. Lines were generated using the Redlich–Kister equation (eq 9).

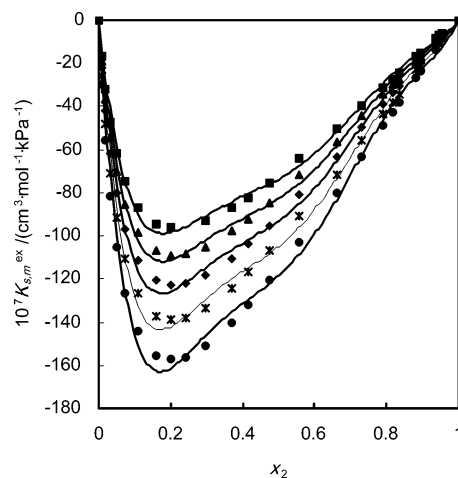


Figure 3. Plot of excess molar isentropic compression $K_{s,m}^{\text{ex}} / (\text{cm}^3 \cdot \text{mol}^{-1} \cdot \text{kPa}^{-1})$ for methanol (1) + PEG400 (2) system against mole fraction of polymer x_2 at different temperatures: \blacksquare , 288.15 K; \blacktriangle , 298.15 K; \blacklozenge , 308.15 K; \times , 318.15 K; \bullet , 328.15 K. Lines were generated using the Redlich–Kister equation (eq 9).

atom of the hydroxyl group of methanol and the oxygen atoms of the poly(ethylene glycol) as well as the packing effect are responsible for the negative V_m^{ex} values obtained for the PEG + methanol system. The excess molar volume values reported for methanol + triethylene glycol¹⁷ and methanol + poly(ethylene glycol) dimethyl ether 250¹⁸ are also negative at different temperatures, and in these systems V_m^{ex} values also become more negative when temperature increases. This behavior of variation of V_m^{ex} with temperature is similar to that of PEG 400 + methanol studied in this work. Valtz et al.¹⁷ reported that such behavior may be explained by the packing effects that become more dominant and increase with temperature.

The negative behavior reflected in Figure 3 for the excess molar isentropic compression for the whole range of composition and different temperatures implies a great difficulty to compress the PEG 400 + methanol solutions than the ideal behavior at different polymer mole fractions and the working temperatures. Figure 3 also shows that as the temperature increases, deviation of methanol + PEG400 solutions from the ideal behavior become larger.

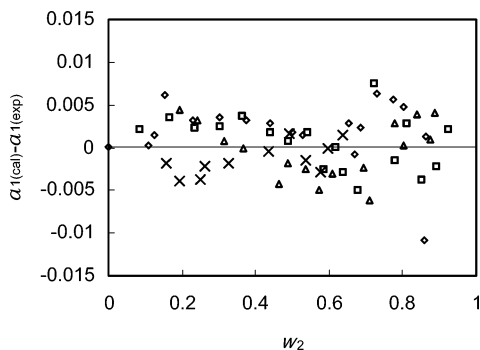


Figure 4. Difference between the experimental and calculated methanol activities with the modified Flory–Huggins model (eq 7) at different temperatures: \diamond , 298.15 K; \square , 308.15 K; \triangle , 318.15 K; \times , 328.15 K.

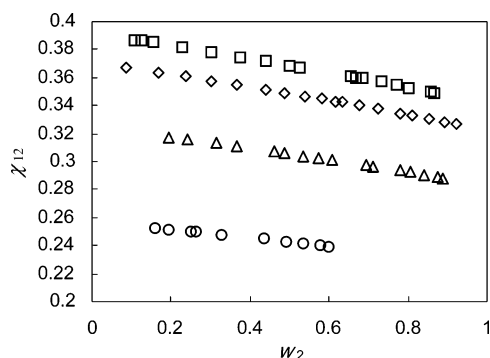


Figure 5. Variation of the Flory parameter χ_{12} with concentration and temperature for methanol (1) + PEG400 (2) system: \square , 298.15 K; \diamond , 308.15 K; \triangle , 318.15 K; \circ , 328.15 K.

Table 4. Flory Parameters Calculated from Equation 6 along with the Standard Deviations for Methanol Activity $\sigma(a_1)$ for the System Methanol (1) + PEG400 (2) at Different Temperatures

T/K	χ_{12}	$\sigma(a_1)$
298.15	0.3522	0.0036
308.15	0.3290	0.0033
318.15	0.2815	0.0046
328.15	0.2329	0.0024

Table 5. Parameters of Equation 9 for V_m^{ex} along with the Corresponding Standard Deviations, $\sigma(V_m^{\text{ex}})$ for Methanol (1) + PEG400 (2) System at Different Temperatures

T	B_0	B_1	B_2	B_3	B_4	$\sigma(V_m^{\text{ex}})$
K						$\text{cm}^3 \cdot \text{mol}^{-1}$
288.15	-3.707	-2.630	-4.229	-4.339	-	0.04
	-3.900	-2.491	-1.262	-4.525	-4.713	0.02
298.15	-3.832	-2.642	-4.405	-4.844	-	0.04
	-4.042	-2.606	-1.363	-4.902	-4.909	0.02
308.15	-4.011	-2.768	-4.712	-5.171	-	0.04
	-4.236	-2.729	-1.452	-5.233	-5.260	0.03
318.15	-4.206	-2.917	-5.021	-5.558	-	0.05
	-4.446	-2.875	-1.545	-5.624	-5.609	0.03
328.15	-4.434	-3.086	-5.405	-5.911	-	0.05
	-4.700	-3.040	-1.552	-5.984	-6.216	0.03

-7394.856, -25.050, -0.147, and 0.0037. To see the performance of eq 7 in representing solvent activity data, the lines generated for methanol activity at (298.15 and 328.15) K were obtained using the parameters of eq 7, and these are also shown in Figure 1. The difference between the experimental and calculated methanol activities with the modified Flory–Huggins model (eq 7) are shown at different temperatures in Figure 4.

Table 6. Parameters of Equation 9 for $K_{s,m}^{\text{ex}}$ along with the Corresponding Standard Deviations $\sigma(K_{s,m}^{\text{ex}})$ for Methanol (1) + PEG400 (2) System at Different Temperatures

T	B_0	B_1	B_2	B_3	B_4	$10^7 \sigma(K_{s,m}^{\text{ex}})$
K						$\text{cm}^3 \cdot \text{mol}^{-1} \cdot \text{kPa}^{-1}$
288.15	-273.96	205.18	-414.22	474.08	-	4.13
	-294.94	190.10	-92.32	494.24	-511.38	2.35
298.15	-305.72	217.62	-465.52	563.73	-	4.82
	-330.66	213.30	-104.16	570.57	-583.10	2.71
308.15	-343.28	243.86	-532.16	649.67	-	5.61
	-372.24	238.85	-112.58	657.61	-677.05	3.17
318.15	-385.90	273.21	-611.30	747.59	-	6.58
	-419.85	267.33	-119.41	756.90	-793.74	3.71
328.15	-434.33	306.18	-704.84	862.53	-	7.74
	-474.34	299.24	-125.07	873.51	-935.54	4.35

In the modified Flory–Huggins model, the Flory parameter χ_{12} is defined as

$$\chi_{12}(T, z) = \frac{d_0}{1 - f(1 - z)} + \frac{d_1}{T[1 - f(1 - z)]} + \frac{d_2 \ln T}{1 - f(1 - z)} \quad (8)$$

The variations of the Flory parameter χ_{12} with PEG mass fraction and the temperature were obtained from eq 8, and the corresponding plot is shown in Figure 5. Figure 5 shows that the modified Flory–Huggins model gives the χ_{12} values, which decrease by increase in temperature, similar to the trend obtained by the original Flory–Huggins model as reported in Table 4.

The excess molar volumes and the excess molar isentropic compressions were correlated by means of the Redlich–Kister equation:²¹

$$\Delta Q = x_2(1 - x_2) \sum_{p=0}^N B_p (2x_2 - 1)^p \quad (9)$$

where B_p represents the fitting coefficients and N is the degree of the polynomial expansion. The standard deviations, σ , between the calculated, ΔQ_{calc} , and the experimental, ΔQ_{exptl} , values have been estimated by using

$$\sigma = \left(\frac{\sum_{i=1}^{n_{\text{DAT}}} (\Delta Q_{\text{exptl}} - \Delta Q_{\text{calc}})^2}{n_{\text{DAT}}} \right)^{1/2} \quad (10)$$

where n_{DAT} is the number of experimental points. The obtained adjustable parameters, B_p , are summarized in Tables 5 and 6 together with the standard deviations, σ . Using eq 9 with $N = 4$ and the corresponding parameters for V_m^{ex} and $K_{s,m}^{\text{ex}}$ reported

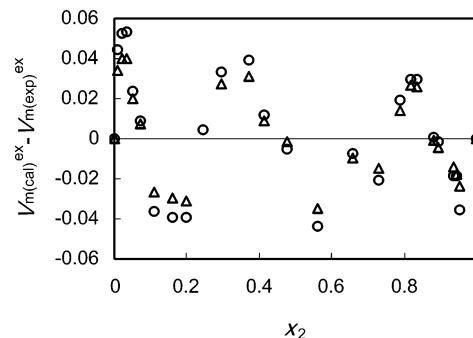


Figure 6. Difference between the experimental and calculated excess molar volumes V_m^{ex} , with the Redlich–Kister equation (eq 9) at two temperatures: \triangle , 288.15 K; \circ , 328.15 K.

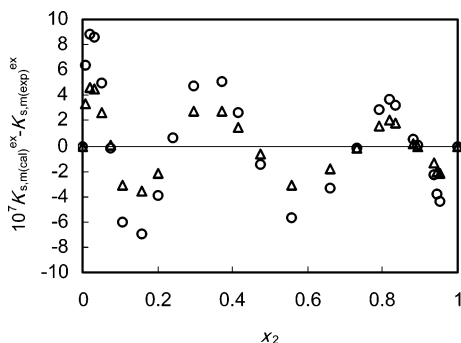


Figure 7. Difference between the experimental and calculated excess molar isentropic compression $K_{s,m}^{ex}$, with the Redlich–Kister equation (eq 9) at two temperatures: Δ , 288.15 K; \circ , 328.15 K.

in Tables 5 and 6, the lines generated for V_m^{ex} and $K_{s,m}^{ex}$ were obtained, and these were also respectively shown in Figures 2 and 3 for the five studied isotherms. Slightly higher standard deviations for V_m^{ex} and $K_{s,m}^{ex}$ are obtained when we use eq 9 with $N = 3$, as can be seen in Tables 5 and 6. The difference between the experimental and calculated excess molar volumes and excess molar isentropic compressions with the Redlich–Kister equation (eq 9) are shown respectively in Figures 6 and 7 at two temperatures (288.15 and 328.15) K as examples. From the $K_{s,m}^{ex}$ values calculated from eq 9, one can easily evaluate the values for isentropic compressibility and speed of sound at a desired concentration with reasonable accuracy.

Conclusions

Accurate methanol activity data were determined for the methanol + PEG400 system at $T = (298.15 \text{ to } 328.15) \text{ K}$, and these data were fitted to the Flory–Huggins model. It was found that both the methanol activity and the Flory parameter decrease as the temperature is increased, indicating an increase of interaction between polymer and methanol by increasing the temperature. From the measured density and speed of sound data at $T = (288.15 \text{ to } 328.15) \text{ K}$, the values for excess molar volume and excess molar isentropic compression for methanol + PEG400 were calculated. Both of these quantities are negative at different concentrations and working temperatures. The negative excess molar volume along with the variation of the Flory parameter with temperature were used to obtain some information with regard to the polymer–solvent interactions and the packing effect. The negative excess molar isentropic compression values imply a greater difficulty to compress the methanol + PEG400 solutions than the ideal behavior at different temperatures considered in this work.

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